DESIGN OF ULTRA-HIGH TEMPERATURE CERAMICS FOR OXIDATION RESISTANCE II: THERMODYNAMIC MODELING OF THE HAFNIUM-TANTALUM-OXYGEN SYSTEM

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Abstract

Phase equilibria in the Hf-Ta-O system and $HfO₂-Ta₂O₅$ isoplethal section are modeled to optimize the thermal stability of HfC-TaC ceramics. Thermodynamic models under the CALculation of PHAse Diagrams (CALPHAD) framework are developed in conjunction with first-principles calculations and experimental data from the literature. Density functional theory (DFT) calculations of binary oxide thermodynamic properties agree with previous experimental data and predict cation disorder in $Hf_6Ta_2O_{17}$. Prior models of the Hf-Ta system, revised ones of the Hf-O system, and present ones of the Ta-O system are included in the ternary modeling. The ternary oxide series, $Hf_{(n-5)/2}Ta_2O_n$, is modeled as three solid solutions with disordered cation sublattices that increase in stability with structure size due to entropic stabilization. The current models predict that $Hf_4Ta_2O_{13}$ is metastable, consistent with lack of experimental evidence supporting its stability. The computed phase diagrams improve upon previous ones and predict that optimal thermal stability occurs at compositions between 3 HfC : 1 TaC and 4 HfC : 1 TaC.

1. Introduction

Ultra-high temperature ceramics (UHTCs) have the highest achievable melting temperatures above $3000 \, \text{°C}$ [1]. Due to the high thermal conductivity of UHTCs, they are considered potential materials for aerospace applications by the National Aeronautics and Space Administration (NASA) [2-4]. In particular, hafnium carbide (HfC) and tantalum carbide (TaC) are the UHTCs with the highest respective melting temperatures of 3900 $\rm{^{\circ}C}$ and 3880 $\rm{^{\circ}C}$ [5]. Nevertheless, HfC and TaC respectively oxidize to form $HfO₂$ and Ta₂O₅ with inadequate protective capability, and severe oxidation within minutes of oxygen exposure above 1600 \degree C limits the application of the carbides. HfC-TaC solution phases have been observed to form a densely packed, thermal-shock resistant, and adherent $Hf_6Ta_2O_{17}$ superstructure that is part of the $Hf_{(n-5)/2}Ta_2O_n$ homologous series [6], where $13 \le n \le 21$ [7]. Hf₆Ta₂O₁₇ has been found to have slower oxygen conduction compared to that of $HfO₂$ and Ta₂O₅ [8, 9], which is desirable for extreme condition applications including hypersonic aircraft design, thermal protection systems, and propulsion systems. However, the phase equilibria in the Hf-Ta-O system and $HfO_2-Ta_2O_5$ isoplethal section, including the stabilities of $Hf_{(n-5)/2}Ta_2O_n$ and

 $Hf_6Ta_2O_{17}$, remain unclear. The present study, described in more detail in an upcoming publication [10], is a continuation of [11] and models the phase equilibria in the ternary system using the CALculation of PHAse Diagrams (CALPHAD) method [12], incorporating the Ta-O models described in previous papers by the authors [11, 13]. The method combines calculated thermochemical properties with experimental data to determine optimized phase diagrams.

2. Review of Prior Experimental Data

2.1. Phases and Transformations in the Hf-Ta-O System.

The ambient pressure, equilibrium phases in the Hf-Ta-O system include hexagonal close-packed (HCP) α -Hf, body-centered cubic (BCC) β-Hf, BCC Ta, three $HfO₂$ polymorphs (monoclinic (M) , tetragonal (T) , and face-centered cubic (C)), two $Ta₂O₅$ polymorphs (orthorhombic (O) and tetragonal (T)), O-H $f_{(n-5)/2}Ta_2O_n$, liquid, and gas [14]. While the stable O-Ta₂O₅ crystal structure is still debated, the β modification refined by Hummel et al. [15] is used for the present density functional theory (DFT) calculations and modeling due to its similarity to $T-Ta₂O₅$ [16], as the authors previously described in detail [11, 13]. The chosen β modification is nonstoichiometric with a

composition of $Ta_{24}O_{62}$ due to partial occupancy of oxygen sites. Crystal structures and phase transformations in the Hf-Ta-O system modeled in the current study are tabulated in the Part I paper [11].

2.2. Cation Disorder and Kinetic Metastability in $O-Hf_{(n-5)/2}Ta_2O_n$

Cation disorder in the O-H $f_{(n-5)/2}Ta_2O_n$ series is very likely based on experimental evidence of it in isostructural O-Zr₆Nb₂O₁₇ [17], O-Zr₁₀Nb₂O₂₅ [18], $O-Zr_6Ta_2O_{17}$ [7], and $O-Hf_6Nb_2O_{17}$ [7]. Recent high temperature oxide melt solution calorimetry experiments of $O-A_6B_2O_{17}$, where A is Hf or Zr and B is Ta or Nb, by Voskanyan et al. [19] provide more direct evidence of cation disorder in O-Hf₆Ta₂O₁₇ and by extension, O-Hf_{(n-5)/2}Ta₂O_n. The study found that each ternary oxide has a positive enthalpy of formation from the respective binary oxides and requires considerable entropy to be stabilized. Since configurational entropy has a much higher magnitude than that of vibrational entropy in solid phases, Voskanyan et al. concluded that the former, arising from cation disorder, stabilizes the structures. Spurling et al. [20] further validated configurational cation disorder in $O-A_6B_2O_{17}$ structures using ex-situ and in-situ high temperature X-ray diffraction (XRD) of ternary and quinary $O-A_6B_2O_{17}$ phases. The exsitu and in-situ experiments respectively resulted in O-H f_6 Ta₂O₁₇ stabilization temperatures of 1273 K and 1423 K, the difference of which with time indicating kinetic metastability. The longer timescale ex-situ value was used in the present study due to its improved accuracy for kinetically metastable O-H $f_6Ta_2O_{17}$. The current CALPHAD models incorporate the cation disorder, as well as kinetic metastability, of $O-Hf_6Ta_2O_{17}$ and by extension, of O-H $f_{(n-5)/2}Ta_2O_n$.

2.3. Prior Studies of the $HfO_2-Ta_2O_5$ Isoplethal Section

The first phase diagram of the HfO_2 -Ta₂O₅ isoplethal section across the entire compositional range was constructed by Turcotte [21] using phase data from thermal arrest measurements, XRD, and electron microscopy. Turcotte concluded that the Ta_2O_5 phase acts as a solid solution or series of homologous structures based on similar XRD patterns between 3.9 and 9 mol% $HfO₂$. A comprehensive $HfO₂-Ta₂O₅$ isoplethal section was constructed by McCormack et al. [14] using thermal arrest measurements and in-situ Xray powder diffraction (XRPD) with a conical nozzle levitator system to accurately measure the invariant transformation, solidus, and liquidus temperatures, equilibrium phases, and reversibility of phase transitions. The phase diagram was revised [22] to satisfy Hollman's Rule [23-25] and was used as the basis for the current CALPHAD models. Inconsistencies in the phase diagram, including lack of entropic stabilization and kinetic metastability and absence of an increase in stability with higher cation disorder [19] of the O- $Hf_{(n-5)/2}Ta_2O_n$ phase, were resolved by the present CALPHAD models.

3. Methods

3.1. First-Principles Calculations

As inputs for the CALPHAD models, 0-K and finite-temperature properties of M-HfO₂, O-Ta₂O₅, and $O-Hf_6Ta_2O_{17}$ were calculated using DFT and phonon calculations under the quasiharmonic approximation (QHA). The ordered form of O- $Hf_6Ta_2O_{17}$, although inconsistent with the 73-90 % cation disorder observed in the present mixing energies and prior experimental studies [19, 20], was used for the 0-K formation enthalpy and phonon calculations to reach manageable computational costs. Figure 1 shows the crystal structures of the oxides used for the first-principles calculations [7, 15, 26]. The DFT calculations run in the Vienna Ab initio Simulation Package (VASP) will be described in detail in an upcoming publication [10]. The local density approximation (LDA) functional [27] was used based on computations of the oxide formation enthalpies using five exchange-correlation functionals. Additional calculated properties include entropies (S) , heat capacities (C_n) , and Gibbs energies (G) of M-HfO₂ and O-Hf₆Ta₂O₁₇, and mixing enthalpies (ΔH_{mix}) of Ta and Hf in O-Ta₂O₅ and $O-Hf_6Ta_2O_{17}$. For validation, the heat capacities were compared to prior data from the literature. Phonon calculations were not performed on the O- $Ta₂O₅$ phase due to the overly high computational costs of modeling the oxygen partial occupancy in the chosen β modification [15].

Figure 1. Crystal structures of $O-Hf_6Ta_2O_{17}$ [7], M-HfO₂ [26], and O-Ta₂O₅ [15] used for the current DFT computations.

0-K formation enthalpies of the oxides were calculated using Equations 1-3, where (E_0) is the total energy of the corresponding structure from full relaxation with DFT. The O_2 energy is that of an isolated molecule in a cubic 10 x 10 x 10 Å cell with spin polarization due to the unpaired electrons in its triplet ground state [28].

$$
\Delta H_{f,0-K}^{O-Ta_2O_5} = E_0(Ta_2O_5) - 2E_0(Ta) - \frac{5}{2}E_0(O_2)
$$
\n(1)

$$
\Delta H_{f,0-K}^{M-HfO_2} = E_0(HfO_2) - E_0(Hf) - E_0(O_2) \quad (2)
$$

$$
\Delta H_{f,0-K}^{O-Hf_6Ta_2O_{17}} = E_0(Hf_6Ta_2O_{17}) - 6E_0(HfO_2) - E_0(Ta_2O_5)
$$
\n(3)

To determine whether the error in the DFT formation enthalpies resulted from the choice of $Ta₂O₅$ structure, further values were calculated for O-Hf₆Ta₂O₁₇ and isostructural O-Hf₆Nb₂O₁₇ [7]. Total energies of the β (*C2mm*), T (*Pmm2*), B (*C2/c*), β (*Pccm*), δ (*P6/mmm*), β (*Pmma*), λ (*Pbam*), *Pmmn*, and γ (*I41/amd*) Ta2O⁵ modifications described in the Part I paper [11], and the P (*I41/amd*) [29, 30], R (*A2/m*) [31], N (*C2/m*) [32], M (*I4/mmm*) [33], and B (*C2/c*) [34]

Nb2O⁵ modifications were input into Equation 3 for the respective formation enthalpy calculations.

Phonon dispersions and vibrational entropies were computed in Phonopy [35] under the harmonic approximation. Heat capacities and Gibbs energies as functions of temperature were calculated in Phonopy using the third-order Birch-Murnaghan [36] equation of state and quasiharmonic approximation (OHA). ΔG and ΔS were used to compute the 298-K enthalpies of formation using Equation 4 [37], in which F is the Helmholtz free energy, E_0 is the enthalpy at 0 K, and F_{vib} is the vibrational free energy contribution.

$$
F(V,T) = E_0(V) + F_{vib}(V,T)
$$
 (4)

Dilute enthalpies of mixing of Ta and Hf in O- Ta_2O_5 were computed by DFT calculations of the $Hf_{23}Ta_1O_{62}$ ($x_{Hf2O5} = 0.96$, $x_{Ta2O5} = 0.04$) and $Hf_1Ta_{23}O_{62}$ ($x_{Hf_2O5} = 0.04$, $x_{Ta_{2}O5} = 0.96$) structures, and those of the $\text{Hf}_{24}\text{O}_{62}$ and $\text{Ta}_{24}\text{O}_{62}$ reference ones, as shown in Equations 5 and 6, respectively.

$$
\Delta H_{mix,dilute}^{Hf_{23}Ta_1O_{62}} = E_0(Hf_{23}Ta_1O_{62}) -
$$

$$
\frac{23}{24}E_0(Hf_{24}O_{62}) - \frac{1}{24}E_0(Ta_{24}O_{62})
$$
 (5)

$$
\Delta H_{mix,dilute}^{Hf_1T a_{23}O_{62}} = E_0(Hf_1T a_{23}O_{62}) -\frac{23}{24}E_0(T a_{24}O_{62}) - \frac{1}{24}E_0(Hf_{24}O_{62})
$$
\n(6)

Enthalpies of mixing of Ta and Hf in O-H f_6 Ta₂O₁₇ between the compositions Hf_8O_{17} and Ta_8O_{17} were calculated using special quasirandom structures (SQS) created using the Alloy Theoretic Automated Toolkit (ATAT) [38]. Equation 7, in which n is an integer from 1 to 3, was used to compute the mixing enthalpies. Mechanical instabilities in the SQS were accounted for using the inflection-detection method [39].

$$
\Delta H_{mix}^{Hf_{8n}Ta_{32-8n}O_{68}} = E_0(Hf_{8n}Ta_{32-8n}O_{68}) - nE_0(Hf_8O_{17}) - (4-n)E_0(Ta_8O_{17}) \tag{7}
$$

3.2. Thermodynamic Models

All phases in the Hf-Ta-O system were modeled using the compound energy formalism (CEF), the details of which are described in the Part I paper [11]. Due to the stoichiometry of the Ta₂O₅ and O- $Hf_{(n-5)/2}Ta_2O_n$ phases [13, 22], their vacancy concentrations and binary excess interaction terms with vacancies, were considered negligible. Binary interaction parameters $({}^{\nu}L_{ij})$ with a linear temperature dependence were added as needed to fit the experimental phase data. Gibbs energy equation parameters were optimized in the Thermo-Calc software [40]. Parameters from the Scientific Group Thermodata Europe (SGTE) database [41] were used for the elemental reference states.

Modeled phases include those in the binary Hf-O, Hf-Ta, and Ta-O systems, as well as O- and T-Ta₂O₅ solid solutions, $Hf_4Ta_2O_{13}$, $Hf_6Ta_2O_{17}$, $Hf_8Ta_2O_{21}$, and a liquid phase. Existing CALPHAD models of the Hf-O system by Wang [42] were modified to correct errors and more closely estimate the maximum oxygen deficiency in the C-HfO² phase, which was previously determined to be 64 at%, by metallographic and chemical analyses [43-45]. Models of the Hf-Ta and Ta-O systems were respectively adopted from those of Guillermet [46] and those published by the authors [11, 13]. The homologous series of O- Ta_2O_5 and T-Ta₂O₅ [21, 22] were modeled as ionic solutions with the sublattice model (Hf^{+4}, Ta^{+5}) ₂(O⁻², Va)₅. Gibbs energy parameters of both phases were optimized to fit experimental phase data [14, 20, 21]. The $Hf_{(n-5)/2}Ta_2O_n$ homologous series [6, 14] was modeled as three solid solutions with kinetic metastability and cation disorder [7, 19, 20] using the sublattice models $(Hf^{+4}, Ta^{+5})_6(O^{-2}, Va)_{13}$, $(Hf^{+4},Ta^{+5})_8(O^2, Va)_{17}$, and $(Hf^{+4},Ta^{+5})_{10}(O^2, Va)_{21}$ to respectively describe O-Hf₄Ta₂O₁₃, O-Hf₆Ta₂O₁₇, and O-Hf₈Ta₂O₂₁. The Gibbs energy parameters of $O-Hf_6Ta_2O_{17}$ were obtained by fitting the SQS mixing enthalpies and from the experimental phase data [14, 20]. Those of O-H f_4 Ta₂O₁₃ and O- $Hf_8Ta_2O_{21}$ were qualitatively adjusted to result in an increase in stability of $Hf_{(n-5)/2}Ta_2O_n$ with larger structure size to reflect cation entropy stabilization

[19]. The liquid phase was modeled by the ionic two-sublattice model developed by Hillert et al. [47], as:

 $(Hf^{+4}, Ta^{+5})_{2y_{0^{-2}}+Qy_{Va^{-Q}}} (O^{-2}, Va^{-Q})_{4y_{Hf^{+4}}+5y_{Ta^{+5}}},$ where $Q = 4y_{Hf^{+4}} + 5y_{Ta^{+5}}$. Gibbs energy parameters were fit to experimental phase data from McCormack et al. [14] and Turcotte [21]. The optimized thermodynamic models of the HfTa-O system will be included in an upcoming publication [10].

4. Results and Discussion

The 298-K DFT formation enthalpy of M -HfO₂ (-1162.93 kJ/mol) using the LDA functional results in a -4.24 % difference from the experimental value in the SGTE database [48] and the 0-K value of the T modification of $O-Ta₂O₅$ $(-2047.38 \text{ kJ/mol})$ results in a $+0.08$ % difference. Despite the high accuracy for the binary oxides, the calculated 0-K values for both $O-Hf_6Ta_2O_{17}$ and O-Hf6Nb2O¹⁷ shown in Table 1 are all negative and significantly differ from the positive values from oxide melt solution calorimetry [19], indicating that the error is not from the choice of Ta or Nb oxide modification. Rather, the error can be attributed to that from the LDA functional as will be shown by the mixing enthalpy results.

Table 1. 0-K DFT formation enthalpies of O- $Hf_6Ta_2O_{17}$ and O-H $f_6Nb_2O_{17}$ from M-HfO₂ [26] and Ta_2O_5 and Nb_2O_5 modifications with available structural data compared to oxide melt solution calorimetric values of 42.94 ± 7.03 kJ/mol and 38.44 \pm 6.75 kJ/mol, respectively [19].

The calculated heat capacity of $M-HfO₂$ under the QHA as a function of temperature is shown in Figure 2 and agrees with that in the SGTE database [48], that calculated using the stoichiometric Hf-O CALPHAD model [42], and

that from prior ab-initio molecular dynamics (AIMD) and QHA calculations [49]. Figure 3 shows the heat capacity of $O-Hf_6Ta_2O_{17}$ calculated under the QHA, which agrees with the current thermodynamic models, as well as with values measured by Li et al. [50] using the comparison method with laser flash analysis. The aforementioned values underestimate those of Perepezko [51], which were obtained using the comparison method with laser flash analysis. The values of Perepezko [51] are expected to be less accurate, as they indicate a sharp rise in heat capacity of a densely packed structure.

Figure 2. Constant pressure heat capacities of M- $HfO₂$ as a function of temperature calculated by the QHA compared to prior AIMD and QHA calculations [49], a stoichiometric CALPHAD model [42], and experimental values from the SGTE database [48].

Figure 3. Constant pressure heat capacities of O- $Hf_6Ta_2O_{17}$ as a function of temperature calculated by the QHA compared to those from the present CALPHAD model and experimental measurements by Li et al. [50] and Perepezko [51].

Mixing enthalpies of Ta and Hf in the $O-Ta_2O_5$ solid solution fitted to phase data of the HfO_{2} - $Ta₂O₅$ system under the current model parameters are shown in Figure 4. SQS mixing enthalpies of Ta and Hf in O-H f_6 Ta₂O₁₇ fitted to a zero- and first-order interaction parameter Redlich-Kister polynomial [52] are shown in Figure 5. Based on the minimum mixing energy at 75 at% Hf, the calorimetric formation enthalpy of 42.94 ± 7.03 kJ/mol [19], and a cationic configurational entropy of 28.05 J/mol/K, a minimum stabilization temperature of 1531 \pm 251 K is predicted for O- $Hf_6Ta_2O_{17}$, consistent with the 1273 K value measured by Spurling et al. [20]. The SQS mixing enthalpies, like the formation enthalpies, of O- $Hf_6Ta_2O_{17}$ are negative in contrast to prior calorimetric findings [19]. Calculated 0-K formation enthalpies of the 75 % cation disorder SQS structure using each modification of $O-Ta₂O₅$ are more negative and have a larger error relative to the positive calorimetric formation enthalpy [19] than that of the ordered structure. This indicates that the DFT error in the signs of the O-H $f_6Ta_2O_{17}$ formation and mixing enthalpies is due to underestimation of the total energies by the LDA functional rather than the choice of $O-Ta₂O₅$ modification or use of ordered ternary oxide structure for the former.

Figure 4. Fit of mixing enthalpies of Hf and Ta in the O-Ta₂O₅ solid solution as a function of Hf content using the ionic solution model based on experimental $HfO_2-Ta_2O_5$ isoplethal section phase data [14]. Dilute mixing enthalpies from DFT calculations are also shown.

Figure 5. Fit of SQS mixing enthalpies of Hf and Ta in O-H f_6 Ta₂O₁₇ as a function of Hf content using nonlinear regression.

The revised Hf-O phase diagram (Figure 6) agrees with experimental phase data in the literature [43- 45, 53]. Hf-Ta-O models optimized using zeroand first-order binary interaction parameters sufficiently describe the experimental $HfO_2-Ta_2O_5$ isoplethal section and will be tabulated in an upcoming publication [10]. A calculated isothermal section of the Hf-Ta-O system at 1773 K is displayed in Figure 7 and better agrees with experimental phase data [14] relative to previous ones [54, 55]. **Prior Revised**

Figure 6. Revised Hf-O phase diagram.

Figure 7. Calculated isothermal Hf-Ta-O section at 1773 K.

The calculated $HfO_2-Ta_2O_5$ isoplethal section between 273 and 3500 K (Figure 8) agrees well with experimental phase data from Spurling et al. [20] and Turcotte [21] but differs from the isoplethal section by McCormack et al. [22]. The phase transitions and invariant equilibria of the latter are compared to current results in Figure 8. Primary differences are those in the stabilities of the O-H $f_{(n-5)/2}Ta_2O_n$ phases and those in the Ta₂O₅rich liquidus. In contrast to the phase diagram of McCormack et al., the current study finds that O- $Hf_4Ta_2O_{13}$, O-H $f_6Ta_2O_{17}$, and O-H $f_8Ta_2O_{21}$ respectively increase in stability due to increasing entropic stabilization and have minimum stabilization temperatures above 298 K due to kinetic metastability. This finding is consistent with more recent experimental ones by Spurling et al. [20] and Voskanyan et al. [19]. In the present study, $Hf_4Ta_2O_{13}$ requires the highest enthalpy of formation and most negative entropy of formation of the ternary oxides for it to be stable, unlike the expected trends $\Delta H_f^{O-Hf_4Ta_2O_{13}} < \Delta H_f^{O-Hf_6Ta_2O_{17}} <$ ∆H $_{f}^{O-Hf_8Ta_2O_{21}}$ and (most negative) $S_{conf}^{Hf_8Ta_2O_{21}}$ < $S_{conf}^{Hf_6Ta_2O_{17}} < S_{conf}^{Hf_4Ta_2O_{13}}$. Due to these discrepancies, as well as lack of data supporting the stability of $A_4B_2O_{13}$ (where A is Hf or Zr and B is Ta or Nb) and the prior concluded metastability of $Hf_3Ta_2O_{11}$ [56], $Hf_4Ta_2O_{13}$ is also concluded to be metastable. Regarding the differences in the Ta_2O_5 -rich liquidus, the current models overestimate the T- $Ta₂O₅$ liquidus relative to that of McCormack et al. In addition, they do not consider a monotectic or miscibility gap due to lack of confirmation of its existence by X-ray scattering data [14], and instead predict a congruent melting point at $x_{HfO2} = 0.145$ and T = 2236 K.

Figure 8. (Left) Calculated HfO₂-Ta₂O₅ isoplethal section. Experimental data from [20] (Green), [14] (Red), and [21] (Blue). The dotted phase boundary of $Hf_4Ta_2O_{13}$ in the calculated isoplethal section indicates metastability. (Right) Experimental HfO₂-Ta₂O₅ isoplethal section constructed and revised by McCormack et al. [14, 22]. Dashed lines specify regions for which additional experiments are needed for validation. Figure from S. J. McCormack et al., "Reply to comments: "In-situ determination of the HfO2‐Ta2O5‐temperature phase diagram up to 3000°C"," *J. Am. Ceram. Soc.,* vol. 102, no. 11, 2019 republished with permission from Elsevier; Permission was conveyed through Copyright Clearance Center, Inc.

5. Conclusions

Thermodynamic properties of $M-HfO₂$ and O- $Ta₂O₅$ are accurately described by DFT calculations and support cation disorder in the O- $Hf_{(n-5)/2}Ta_2O_n$ phase in agreement with prior experimental findings. The modified phase diagram of the Hf-O system and new ones of the Hf-Ta-O system agree with experimental phase data, and the latter show that the O-Hf $_{(n-5)/2}Ta_2O_n$ phase increases in stability with higher $HfO₂$ content and has minimum stabilization temperatures above 298 K. These phenomena occur due to cation entropy stabilization and kinetic metastability, respectively. Consequently, HfC-TaC UHTCs with compositions between 3 HfC : 1 TaC and 4 HfC : 1 TaC are predicted to have the optimal thermal stability. Further studies pertaining to the thermodynamics and kinetics of oxidation are required to compare oxygen transport in the O-H $f_{(n-5)/2}Ta_2O_n$ phase to those of $HfO₂$ and Ta₂ $O₅$ and are currently being carried out by the authors.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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